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Oscillation of Super-linear second Order Nonlinear Differential Equations with Damping Term

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The study of differential equations has been the object of many researchers over the last decades. Different approaches and various techniques have been adopted to investigate the qualitative properties of their solutions. Recently and driven by their widespread applications, the investigation of differential equations of second order has drawn significant attention. The oscillation of solutions has been the main features that have attracted consideration. Therefore, it has been intended to use the Riccati Transformation Technique for obtaining several new oscillation criteria for different classes of nonlinear differential equations of the second order with a damping term. Oscillatory behavior has taken into account through this study of solutions of some differential equations. Comparisons between our results and the previously known results have presented. The relevance of our theorems has been clear due to carefully selected examples. As a conclusion, our aim is to provide some results to improve and/or extend some of well-known results in the literature.

1 Introduction

This paper concerned with oscillation of the solution to the damped ordinary differential equation of the form:

$$(r(t)\psi(x(t))x'(t))' + h(t)x'(t) + q(t)g(x(t)) = H(t, x(t), x'(t)), \quad (1)$$

where r, ψ, h and q are continuous functions on the interval $[t_0, \infty)$, $t_0 \geq 0$ $r(t)$ is a positive function, ψ is continuous function on the real line \mathbb{R} , with $\psi(x) > 0, \forall x \in \mathbb{R}$ and g is continuously differentiable function on the real line \mathbb{R} except possible at 0 with $xg(x) > 0$ and $g'(x) \geq k > 0$ for all $x \neq 0$. H is a

continuous function on $[t_0, \infty) \times \mathbb{R}^2$ with

$$\frac{H(t, x(t), x'(t))}{g(x(t))} \leq p(t) \text{ for all } t \in [t_0, \infty)$$

Throughout this study, our attention is only to the solutions of the differential equation (1) that exist on

some ray $[t_x, \infty)$ where t_x may depend on the particular solution.

In the past decades, the problems regarding the study of oscillation criteria of differential equations with damping have become an important area of research because such equations arise in many real life problems; see the research papers (Ayanlar & Tiryaki, 2000, Elabbasy & Elhaddad, 2007, Kirane & Rogovchenko, 2001, Mustafa et al., 2004, Nagabuchi & Yamamoto, 1988, Naito, 1974, S. & Rogovchenko, 2003, Rogovchenko & Tuncay, 2007, Rogovchenko & Tuncay, 2008, Rogovchenko & Tuncay, 2009, Saker et al., 2003, Tunc & Avci, 2012, Wang & Song, 2013, Xiaoling & Chenghui, 2013, Zhang & Song, 2011 and Zheng, 2006) and the references cited therein. In particular, second order-damped differential equations are used in the study of vehicle noise, vibration and harshness of vehicles (NVH). In what follows, we present the background details that motivate the contents of this paper.

(Elabbasy *et al.* 2005, Lu & Meng, 2007, and Rogovchenko & Tuncay, 2008), established some new oscillation conditions of Kamanev and Philos type for the Eq. (1) with $\psi(x(t)) = 1, H(t, x(t), x'(t)) = 0$ and they did not sign conditions on $h(t), q(t)$. Rogovchenko & Tuncay, (2009) considered Eq. (1) with $H(t, x(t), x'(t)) = 0$ and Berkan (2008), continued the investigation of new oscillation of Eq. (1) but he put $H(t, x(t), x'(t)) = H(t)$. Zhang & Song, (2011) considered equation (1) when replaced explicit function $q(t)g(x(t))$ by implicit function $Q(t, x)$ they obtained certain necessary criteria for oscillation of Eq. (1).

Wang & Song, (2013), established certain oscillation standards for Eq. (1), with $H(t, x(t), x'(t)) = 0, \frac{g(x(t))}{x(t)} \geq k > 0$ for $x \neq 0$.

Xiaoling & Chenghui, (2013), established an important extension of the celebrated oscillation criteria for (1), they studied it with $\psi(x(t)) = 1, g(x(t)) = g(x(\tau(t)))$ and $H(t, x(t), x'(t)) = 0$.

In the same way, we localize generalize Reccati technique to derive new oscillation conditions for Eq. (1). Our results are more general than the previous results. Precisely chosen examples are provided to demonstrate the influence of impulses on the oscillatory actions of all solutions in this class.

2 Definitions

Definition 1. A solution $x(t)$ of Eq. (1) is oscillatory if it has arbitrarily large zeros; otherwise, we call it non-oscillatory. The Eq. (1) is oscillatory if all of its solutions oscillate.

Definition 2. Equation (1) is said to be super-linear if

$$0 < \int_{\pm \varepsilon}^{\pm \infty} \frac{du}{g(u)} < \infty \text{ for every all } \varepsilon > 0$$

3 Oscillation Results

In this section, I introduce some theorems that include new conditions for ensuring the oscillation of solutions equation (1).

Theorem 1: Suppose that

- (1) $c_1 \leq \psi(x(t)) \leq c_2$ for all $x \in \mathbb{R}$ and
- (2) $h(t) \leq 0$ for $t \geq t_0$ hold.

Let ρ be a continuously differentiable positive function over $[T, \infty)$ such that $\rho'(t) \geq 0$

over $[T, \infty)$; $(\rho'(t)r(t))' \leq 0$ and such that

$$(3) \lim_{t \rightarrow \infty} \int_{T_0}^t \frac{1}{\rho(s)r(s)} ds = \infty,$$

(4)

$$\lim_{t \rightarrow \infty} \int_{T_0}^{\infty} R(s) ds = \infty; R(s) = \rho(s)[q(s) - p(s)] - \frac{1}{4A} \frac{(\rho'(s))^2}{\rho(s)} r(s),$$

A is constant.

Then, every solution of equation (1) is oscillatory.

Proof: Without loss of generality, we may hypothesize that there exists a solution $x(t)$ of equation (1) such that

$$x(t) > 0 \text{ on } [T, \infty), \text{ for some } T \geq t_0 \geq 0.$$

Define

$$\omega(t) = \frac{r(t)\psi(x(t))x'(t)}{g(x(t))}, t \geq T \tag{1}$$

For all $t \geq T_0$ then differentiating Equality (1) and using Eq. (1), we obtain

$$\left(\frac{r(t)\psi(x(t))x'(t)}{g(x(t))} \right)' \leq \frac{H(t, x'(t), x(t))}{g(x(t))} - q(t) - \frac{h(t)x'(t)}{g(x(t))} - \frac{r(t)\psi(x(t))x'(t)g'(x(t))x'(t)}{g^2(x(t))}$$

Since $g'(x(t)) \geq k$ and using the condition (1) we have

$$\left(\frac{r(t)\psi(x(t))x'(t)}{g(x(t))} \right)' \leq -[q(t) - p(t)] - \frac{h(t)x'(t)}{g(x(t))} - \frac{kr^2(t)\psi^2(x(t))(x'(t))^6}{c_2 g^2(x(t))}, t \geq T \tag{2}$$

Multiplying the inequality (2) by $\rho(t)$ and integrate form T to t we obtain

$$\frac{\rho(t)r(t)\psi(x(t))x'(t)}{g(x(t))} \leq C_T - \int_T^t \rho(s)[q(s) - p(s)] ds - \int_T^t \frac{\rho(s)h(s)x'(s)}{g(x(s))} ds + \int_T^t \left[\rho'(s)\omega(s) - A \frac{\rho(s)}{r(s)} \omega^2(s) \right] ds \tag{3}$$

$$\text{Where } C_T = \frac{\rho(T)r(T)\psi(x(T))x'(T)}{g(x(T))}.$$

By using complement square, the inequality presented by (3) can be written as

$$\frac{\rho(t)r(t)\psi(x(t))x'(t)}{g(x(t))} \leq C_T - \int_T^t \rho(s)[q(s) - p(s)] ds - \int_T^t \frac{\rho(s)h(s)x'(s)}{g(x(s))} ds + \int_T^t \left[A \frac{\rho(s)}{r(s)} \left(W^2(s) - \left(\frac{\rho'(s)r(s)}{2A\rho(s)} \right)^2 \right) \right] ds \tag{4}$$

By the Bonnet theorem, for a fixed $\varepsilon_t \in [T, t]$ such that

$$- \int_T^t \frac{\rho(s)h(s)x'(s)}{g(x(s))} ds = -\rho(T)h(T) \int_T^{\varepsilon_t} \frac{x'(s)}{g(x(s))} ds$$

Since $(-\rho(t)h(t)) \geq 0$ and the equation (1) is super-linear, we have

$$-\infty < \int_T^t -\rho(s)h(s) \frac{x'(s)}{g(x(s))} ds \leq B, \tag{5}$$

where $B = -\rho(T)h(T) \int_{x(T)}^{\infty} \frac{du}{g(u)}$.

By (5) and the condition (4), (3) become

$$\frac{\rho(t)r(t)\psi(x(t))x'(t)}{g(x(t))} \leq C_T + B_1 - \int_T^t R(s) ds$$

By the condition (4), we have

$$\lim_{t \rightarrow \infty} \frac{\rho(t)r(t)\psi(x(t))x'(t)}{g(x(t))} = -\infty.$$

Thus, there exists $T_1 \geq T$ such that $x'(t) < 0$ for all $t \geq T_1$. The condition (4) also implies that there exists $T_2 \geq T_1$ such that

$$\int_{T_1}^{T_2} \rho(s)(q(s) - p(s)) ds = 0 \text{ and } \int_{T_2}^t \rho(s)(q(s) - p(s)) ds \geq 0 \text{ for } t \geq T_2$$

Multiplying equation (1) by $\rho(t)$, from the definitions of the functions and condition (2), we get

$$\rho(t)(r(t)\psi(x(t))x'(t))' + \rho(t)h(t)x'(t) + \rho(t)q(t)g(x(t)) = \rho(t)H(t, x(t), x'(t)) \rho(t)(r(t)\psi(x(t))x'(t))' + \rho(t)g(x(t))q(t)$$

$$\leq \rho(t)g(x(t))p(t), t \geq T_2. \tag{6}$$

Integrate the inequality (6) from T_2 to t we obtain

$$\rho(t)r(t)\psi(x(t))x'(t) \leq \rho(T_2)r(T_2)\psi(x(T_2))x'(T_2) + \int_{T_2}^t \rho'(s)r(s)\psi(x(s))x'(s) ds - g(x(t)) \int_{T_2}^t \rho(s)(q(s) - p(s)) ds + \int_{T_2}^t g'(x(s))x'(s) \int_{T_2}^s \rho(u)(q(u) - p(u)) du ds$$

By the condition (1) and the Bonnet's theorem, for $t \geq T_2$ there exists $\gamma_t \in [T_2, t]$ such that

$$c_2\rho(t)r(t)x'(t) \leq \rho(T_2)r(T_2)\psi(x(T_2))x'(T_2) + c_1\rho'(T_2)r(T_2)[x(\gamma_t) - x(T_2)] - g(x(t)) \int_{T_2}^t \rho(s)(q(s) - p(s)) ds + \int_{T_2}^t g'(x(s))x'(s) \int_{T_2}^s \rho(u)(q(u) - p(u)) du ds, t \geq T_2$$

Thus

$$c_2\rho(t)r(t)x'(t) \leq \rho(T_2)r(T_2)\psi(x(T_2))x'(T_2), t \geq T_2$$

Dividing the last inequality by $\rho(t)r(t)$, integrate from T_2 to t and the condition (3), we have

$$c_2x(t) \leq c_2x(T_2) + \rho(T_2)r(T_2)\psi(x(T_2))x'(T_2) \int_{T_2}^t \frac{ds}{\rho(s)r(s)} \rightarrow -\infty$$

as $t \rightarrow \infty$, that is a inconsistency to the fact that $x(t) > 0$ for $t \geq T$. This complete the proof.

Theorem 2: Suppose that the condition (1) hold, and

$$(5) \int_T^{\infty} \frac{ds}{r(s)} \leq k_1, k_1 > 0$$

$$(6) \int_{\pm\varepsilon}^{\pm\infty} \frac{\psi(u)du}{g(u)} < \infty \text{ for all } \varepsilon > 0.$$

Furthermore, suppose that there exists a positive continuous differentiable function ρ on the interval $[t_0, \infty)$ with $\rho(t)$ is a non-decreasing function on the interval $[t_0, \infty)$ such that

$$(7) \limsup_{t \rightarrow \infty} \int_T^t \frac{1}{r(s)\rho(s)} \int_T^s \rho(u) [q(u) - p(u) - \frac{h^2(u)}{4c_1kr(u)}] du ds = \infty,$$

where $\rho: [t_0, \infty) \rightarrow (0, \infty)$.

Thus, each solution of super-linear equation (1) is oscillatory.

Proof: Without loss of generality, we can suppose that there exists a solution $x(t)$ of equation (1) such that $x(t) > 0$ on $[T, \infty)$ for some $T \geq t_0 \geq 0$. Define

$$\omega(t) = \frac{\rho(t)r(t)\psi(x(t))x'(t)}{g(x(t))}, t \geq T$$

This and by the condition (1) and Eq. (1), we have

$$\begin{aligned} \omega'(t) \leq & \rho(t)p(t) - \frac{\rho(t)h(t)x'(t)}{g(x(t))} - \rho(t)q(t) \\ & + \frac{\rho'(t)}{\rho(t)}\omega(t) - \frac{c_1k\rho(t)r(t)(x'(t))^2}{g^2(x(t))}, t \geq T \end{aligned}$$

Thus for $t \geq T$, we have

$$\begin{aligned} \rho(t) \left(\frac{\omega(t)}{\rho(t)} \right)' \leq & \rho(t)p(t) - \rho(t)q(t) - \frac{\rho(t)h(t)x'(t)}{g(x(t))} \\ & - \frac{c_1k\rho(t)r(t)(x'(t))^2}{g^2(x(t))}, t \geq T \end{aligned}$$

$$\begin{aligned} \rho(t)[q(t) - p(t)] \leq & -\rho(t) \left(\frac{\omega(t)}{\rho(t)} \right)' \\ & - \frac{\rho(t)h(t)x'(t)}{g(x(t))} \\ & - \frac{c_1k\rho(t)r(t)(x'(t))^2}{g^2(x(t))}, t \geq T \end{aligned}$$

Integrate from T to t , we obtain

$$\begin{aligned} \int_T^t \rho(s)[q(s) - p(s)] ds & \leq \int_T^t -\rho(s) \left(\frac{\omega(s)}{\rho(s)} \right)' ds \\ & - \int_T^t \frac{\rho(s)h(s)x'(s)}{g(x(s))} ds \\ & - \int_T^t \frac{c_1k\rho(s)r(s)(x'(s))^2}{g^2(x(s))} ds, \quad t \geq T \\ & - \int_T^t \frac{\rho(s)h(s)}{g(x(s))} + \frac{c_1k\rho(s)r(s)(x'(s))^2}{g^2(x(s))} ds \\ & = - \int_T^t \left[\sqrt{c_1k\rho(s)r(s)} \frac{x'(s)}{g(x(s))} \right]^2 ds \\ & \quad + \frac{1}{2} \sqrt{\frac{\rho(s)}{c_1kr(s)}} h(s) \Bigg] ds \\ & \quad + \frac{1}{4kc_1} \int_T^t \frac{\rho(s)h^2(s)}{r(s)} ds \\ & \leq \frac{1}{4kc_1} \int_T^t \frac{\rho(s)h^2(s)}{r(s)} ds \end{aligned}$$

By the Bonnet's theorem, since $\rho(t)$ is a non-decreasing function on the interval $[t_0, \infty)$, there exists $T_1 \in [T, t]$ such that

$$\begin{aligned} - \int_T^t \rho(s) \left(\frac{\omega(s)}{\rho(s)} \right)' ds & = -\rho(t) \int_{T_1}^t \left(\frac{\omega(s)}{\rho(s)} \right)' ds \\ & = -\rho(t) \int_{T_1}^t \left(\frac{d}{ds} \frac{\omega(s)}{\rho(s)} \right) ds \\ & = -\rho(t) \frac{\omega(t)}{\rho(t)} + \rho(t) \frac{\omega(T_1)}{\rho(T_1)} \end{aligned} \tag{7}$$

From the inequalities (7) and (5) in the inequality (4), we have

$$\begin{aligned} \int_T^t \rho(s) \left[q(s) - p(s) - \frac{h^2(s)}{4c_1kr(s)} \right] ds \\ \leq -\rho(t) \frac{\omega(t)}{\rho(t)} + \rho(t) \frac{\omega(T_1)}{\rho(T_1)}. \end{aligned}$$

$$\begin{aligned} \int_T^t \rho(s) \left[q(s) - p(s) - \frac{h^2(s)}{4c_1kr(s)} \right] ds \leq & -\omega(t) \\ & + \rho(t) \frac{\omega(T_1)}{\rho(T_1)}. \end{aligned}$$

$$\int_T^t \rho(s) \left[q(s) - p(s) - \frac{h^2(s)}{4c_1kr(s)} \right] ds \leq -\omega(t)$$

$+ \rho(t) \frac{\omega(T_1)}{\rho(T_1)}$. Integrating the last inequality divided by $\rho(t)r(t)$ from T to t , taking the limit superior on both sides and by conditions (5) and (6), we have

$$\begin{aligned} \limsup_{t \rightarrow \infty} \int_T^t \frac{1}{r(s)\rho(s)} \int_T^s \rho(u) \left[q(u) - p(u) - \frac{h^2(s)}{4c_1kr(s)} \right] duds \\ \leq \limsup_{t \rightarrow \infty} \left\{ \left(\frac{\omega(T_1)}{\rho(T_1)} \right) \int_T^t \frac{ds}{r(s)} - \int_{x(T)}^{x(t)} \frac{\psi(u)du}{g(u)} \right\} \\ < \infty \quad \text{as } t \rightarrow \infty, \end{aligned}$$

which contradicts to the condition (7). Hence, the proof is completed.

4 Discussion

A set of new oscillation conditions are stated and proved which extend and improve previous oscillation criteria and cover the cases which are not covered by known results. Some of illustrative examples are provided to show the applications of the oscillation criteria and the comparisons between our results and previous results in the literature.

Example1: Consider the following differential equation

$$\left(\frac{1}{t} \left(\frac{x^6(t) + 2}{x^6(t) + 1}\right) x'(t)\right)' - t^2 x'(t) + \left(t + \frac{\sin t}{t}\right) x^5(t) = \frac{2x^{12}(t) \sin t \cos(x'(t)+1)}{(x^7+1)t^3}, t \geq \frac{\pi}{2}$$

Here,

$$r(t) = \frac{1}{t}, h(t) = -t^2, q(t) = t + \frac{\sin t}{t}, g(x) = x^5, H(t, x(t), x'(t)) = \frac{2x^{12}(t) \sin t \cos(x'+1)}{(x^7+1)t^3},$$

for all $x \neq 0$ and $t > 0$.

$$\frac{H(t,x(t),x'(t))}{g(x(t))} = \frac{2x^{12}(t) \sin t \cos(x'+1)}{(x^7+1)t^3} \times \frac{1}{x^5(t)} \leq \frac{2}{t^3} = p(t)$$

for all $x \neq 0$ and $t > 0$. $\psi(x) = \frac{x^6+2}{x^6+1}$ and

$$1 \leq \psi(x) \leq 2 \text{ for all } x \in \mathbb{R}, \rho(t) = t, \rho'(t)r(t) = \frac{1}{t} > 0, (\rho(t)h(t))' = -3t^2 < 0 \text{ and } (\rho'(t)r(t))' = \left(\frac{1}{t}\right)' = \frac{-1}{t^2} < 0 \text{ for all } t > 0.$$

So, can note that

$$\int_{t_0}^{\infty} \frac{ds}{\rho(s)r(s)} = \int_{t_0}^{\infty} ds = \infty, R(s) = s^2 + \sin s - \frac{2}{s^2} - \frac{1}{4As^2} \int_{t_0}^{\infty} R(s)ds = \infty.$$

All conditions of Theorem 1 are satisfied; thus, the given equation is oscillatory.

Example2: Consider the following differential equation

$$\left(\frac{(x^2(t) + 2)}{t^4(x^2(t) + 1)} x'(t)\right)' + \frac{x'(t)}{t^5} + t^4 x^5(t) = \frac{x^5(t) \cos(x(t))}{t^9}, t > 0.$$

We note that

$$r(t) = \frac{1}{t}, \psi(x) = \frac{x^2(t)+2}{x^2(t)+1} > 0 \text{ and } 1 \leq \psi(x) \leq 2,$$

$$\text{for all } x \in \mathbb{R}, h(t) = \frac{1}{t^5}, \text{ and } \frac{H(t,x(t),x'(t))}{g(x(t))} = \frac{\cos(x(t))}{t^9} \leq \frac{1}{t^9} = p(t) \text{ for all } t > 0 \text{ and } x \neq 0.$$

Let $\rho(t) = t^6$ such that

$$\limsup_{t \rightarrow \infty} \int_T^t \frac{1}{r(s)\rho(s)} \int_T^s \rho(u) \left[q(u) - p(u) - \frac{h^2(u)}{4c_1 k r(u)} \right] du ds = \limsup_{t \rightarrow \infty} \int_T^t \frac{1}{s^2} \int_T^s u^6 \left[u^4 - \frac{1}{u^9} - \frac{1}{4c_1 k u^6} \right] du ds = \limsup_{t \rightarrow \infty} \left[\frac{s^{10}}{110} - \frac{1}{6s^3} + \frac{1}{8c_1 k T s^2} + \left(\frac{T^{11}}{11s} + \frac{1}{2sT^2} + \frac{1}{4c_1 k s T} \right) \right] = \infty.$$

All conditions of Theorem2 are satisfied and hence each solution of the given equation is oscillatory.

Remark1: Theorem1 and Theorem 2 extend and improve results of (Elabbasy & Elzeine, 2011, Remili, 2008, and Results of Xhevair & Elisabeta, 2014).

Remark2: Remili, (2008) has established some oscillation results for Eq. (1) with $\psi(x(t)) = 1, h(t) = 1$, these results required that $r(t) \leq a_1$ and

$$\liminf_{t \rightarrow \infty} \int_T^t Z(s)ds > -\lambda, \quad (\lambda > 0)$$

for all large $T; Z(s) = R(s)[q(s) - p(s)]$, which are not required in Theorem 2.

5 Conclusions

To sum up, a set of new oscillation conditions are stated and proved which extend and improve previous oscillation criteria and cover the cases which are not covered by known results. Further, we introduced some illustrative examples. Remarks were also included to show the evidence of our main results.

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