



Some New Sufficient Conditions for The Oscillation of Nonlinear Ordinary Differential Equations

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Abstract: In this research, we shall discuss the oscillation behavior of solutions for second order nonlinear differential equations. By employing a generalized Riccati transformation and an averaging technique, we derive some new oscillation conditions for all solutions. Our results are extension for some well known oscillation results. The relevance of our results is illustrated by some examples.

Keywords. Sufficient Conditions, Oscillation, Nonlinear Differential Equations.

1. Introduction:

This paper is concerned with the problem of the oscillation of solutions for the following second order nonlinear ordinary differential equation:

$$(r(t)\psi(x(t))\dot{x}(t))' + G(t, x(t)) = H(t, \dot{x}(t), x(t)) \quad , \quad t \geq t_0$$

(1)



Where

$$r \in C(I, \mathbb{R}_+), \psi \in C(\mathbb{R}, \mathbb{R}_+), I = [t_0, \infty), \mathbb{R} = (-\infty, \infty), \mathbb{R}_+ = (0, \infty)$$

Throughout this paper, we shall impose the following conditions :

O_1 : $G \in C(I \times \mathbb{R}, \mathbb{R})$ and there exists $q \in C(I, \mathbb{R})$ such that $\frac{G(t, x(t))}{g(x(t))} \geq q(t)$ and $g \in C(\mathbb{R}, \mathbb{R})$ such that $xg(x) > 0$ and $g'(x) \geq k > 0$ for $x \neq 0$.

O_2 : $H \in C(I \times \mathbb{R}^2, \mathbb{R})$ and there exists $m \in C(I, \mathbb{R}_+)$ such that $\frac{H(t, \dot{x}(t), x(t))}{g(x(t))} \leq m(t)$ for $x \neq 0$.

In the paper of Amhalhil [2], the oscillatory behavior of Equation (1) was discussed in the case when the right hand side equal to zero; that is when $H(t, \dot{x}(t), x(t)) = 0$. The oscillation property has wide applications especially physical sciences, technology and chemical reaction, one can see Agarwal et al. [1], Manojlovic [8] and Tiryaki et al. [10]. However, among numerous papers dealing with this property we choose to refer to Elabbasy [3], Grace et al. [5], and the recent paper of Elabbasy et al. [4] and Salhin [9].



Here, in this paper, it will be assumed that the solutions of Equation (1) exist for $t_0 \geq 0$. Further, a solution $x(t)$ of Equation (1) is called regular (infinity continuable) if $x(t)$ exists for all $t \geq t_0$. A regular solution of Equation (1) is called oscillatory if it has arbitrarily large zeros and non oscillatory if it is eventually positive or eventually negative.

2. Main Results:

Theorem 2.1: Suppose that

$$O_3: r(t) \leq l_1 \quad \forall t \geq t_0,$$

$$O_4: \frac{1}{\psi(x)} \geq l_2 \quad \text{for all } x \in \mathfrak{R},$$

$$O_5: \lim_{t \rightarrow \infty} \frac{1}{t} \int_{t_0}^t \int_{t_0}^s [q(u) - m(u)] du ds = \infty.$$

Then every solution of Equation (1) is oscillatory.

Proof. Without loss of generality, we may assume that there exists a solution $x(t) > 0$ on $[t_0, \infty)$ for some $T_1 \geq t_0 > 0$. Define

$$\omega(t) = \frac{r(t)\psi(x(t))\dot{x}(t)}{g(x(t))}, \quad t \geq T_1. \quad (2)$$



Then, for $t \geq T_1$, we get

$$\dot{\omega}(t) = \frac{(r(t)\psi(x(t))\dot{x}(t))^{\bullet}}{g(x(t))} - \frac{r(t)\psi(x(t))\dot{x}^2(t)g'(x(t))}{g^2(x(t))},$$

And so, from Equation (1), we have

$$\dot{\omega}(t) = \frac{H(t, \dot{x}(t), x(t))}{g(x(t))} - \frac{G(t, x(t))}{g(x(t))} - \frac{r(t)\psi(x(t))\dot{x}^2(t)g'(x(t))}{g^2(x(t))},$$

Then, for every $t \geq T_1$, we obtain

$$\dot{\omega}(t) \leq m(t) - q(t) - \frac{r(t)\psi(x(t))\dot{x}^2(t)g'(x(t))}{g^2(x(t))},$$

(3)

Thus, for every $t \geq T_1$, we have

$$\omega(t) \leq \omega(T_1) - \int_{T_1}^t \frac{r(s)\psi(x(s))\dot{x}^2(s)g'(x(s))}{g^2(x(s))} ds - \int_{T_1}^t [q(s) - m(s)] ds.$$

Then, by definition of ω , we get

$$\frac{r(t)\psi(x(t))\dot{x}(t)}{g(x(t))} \leq C - \int_{T_1}^t \frac{r(s)\psi(x(s))\dot{x}^2(s)g'(x(s))}{g^2(x(s))} ds - \int_{T_1}^t [q(s) - m(s)] ds, \quad t \geq T_1,$$

Where $C = \omega(T_1)$ is a real constant. Now, from conditions O_3 and O_4 , we have

$$\frac{r(t)\psi(x(t))\dot{x}(t)}{g(x(t))} \leq C - \frac{kl_2}{l_1} \int_{T_1}^t \left[\frac{r(s)\psi(x(s))\dot{x}(s)}{g(x(s))} \right]^2 ds - \int_{T_1}^t [q(s) - m(s)] ds,$$

Integrating again and dividing by t , we obtain

$$\frac{1}{t} \int_{T_1}^t \frac{r(s)\psi(x(s))\dot{x}(s)}{g(x(s))} ds \leq C \left(1 - \frac{T_1}{t}\right) - \frac{d}{t} \int_{T_1}^t \int_{T_1}^s \left[\frac{r(u)\psi(x(u))\dot{x}(u)}{g(x(u))} \right]^2 duds - \frac{1}{t} \int_{T_1}^t \int_{T_1}^s [q(u) - m(u)] duds, \quad t \geq T_1$$

where $d = \frac{kl_2}{l_1}$ is a positive constant. Then, for $t \geq T_1$, by condition O_5 , we obtain

$$\lim_{t \rightarrow \infty} \frac{1}{t} \int_{T_1}^t \frac{r(s)\psi(x(s))\dot{x}(s)}{g(x(s))} ds = -\infty \quad (4)$$

Now, defining

$$V(t) = \left| \int_{T_1}^t \frac{r(s)\psi(x(s))\dot{x}(s)}{g(x(s))} ds \right|, \quad (5)$$

and applying Schwarz's inequality, we have

$$V^2(t) = \left| \int_{T_1}^t \frac{r(s)\psi(x(s))\dot{x}(s)}{g(x(s))} ds \right|^2 \leq \int_{T_1}^t \left| \frac{r(s)\psi(x(s))\dot{x}(s)}{g(x(s))} \right|^2 ds * \int_{T_1}^t |1|^2 ds \leq (t - T_1) \int_{T_1}^t \left| \frac{r(s)\psi(x(s))\dot{x}(s)}{g(x(s))} \right|^2 ds, \quad t \geq T_1.$$

Thus,

$$V^2(t) \leq t \int_{T_1}^t \left| \frac{r(s)\psi(x(s))\dot{x}(s)}{g(x(s))} \right|^2 ds.$$

(6)

Using condition O_5 we obtain

$$-\frac{V(t)}{t} + \frac{d}{t} \int_{T_2}^t \int_{T_2}^s \left[\frac{r(u)\psi(x(u))\dot{x}(u)}{g(x(u))} \right]^2 dud s \leq 0.$$

Then, for all $t \geq T_2$, we have

$$-\frac{V(t)}{t} + \frac{d}{t} \int_{T_2}^t \frac{V^2(s)}{s} ds \leq 0,$$

which follows that

$$\frac{d}{t} \int_{T_2}^t \frac{V^2(s)}{s} ds \leq \frac{V(t)}{t}.$$

(7)

Thus, for all $t \geq T_2$, inequality (7) becomes

$$\frac{d^2}{t^2} \left[\int_{T_2}^t \frac{V^2(s)}{s} ds \right]^2 \leq \frac{V^2(t)}{t^2}.$$

Then, for $t \geq T_2$, we define

$$\varphi(t) = \int_{T_2}^t \frac{V^2(s)}{s} ds.$$

(8)

Then, we get

$$\frac{d^2}{t} \leq \frac{\varphi'(t)}{\varphi^2(t)}.$$

Integrating the last inequality from T_2 to t , we obtain

$$d^2 \ln\left(\frac{t}{T_2}\right) \leq \frac{1}{\varphi(T_2)} - \frac{1}{\varphi(t)} \leq \frac{1}{\varphi(T_2)}.$$

(9)

This is a contradiction and hence the proof is complete.

Example 2.1: Consider the equation:

$$\left[\left(\frac{1}{t+1} \right) \left(\frac{x^2(t)+1}{x^2(t)+2} \right) \dot{x}(t) \right]' + x(t)(t^2 + x^2(t)) = \frac{x^3(t)\dot{x}^2(t)}{t^3(\dot{x}^2(t)+1)(x^2(t)+1)}, \quad t > 0.$$

(10)

We note that

$$0 < r(t) = \frac{1}{t+1} < 1 \quad \text{for all } t > 0,$$

$$\psi(x) = \frac{x^2 + 1}{x^2 + 2} > 0 \quad \text{and} \quad \frac{1}{\psi(x)} = \frac{x^2 + 2}{x^2 + 1} > 1 \quad \forall x \in \mathfrak{R},$$

$$\frac{G(t, x(t))}{g(x(t))} = \frac{x(t)(t^2 + x^2(t))}{x(t)} = (t^2 + x^2(t)) \geq t^2 = q(t) \quad \text{for } x \neq 0 \text{ and } t \in [t_0, \infty)$$

$$\text{and } xg(x) = x^2 > 0 \quad \text{and} \quad g'(x) = 1 > 0 \quad \text{for } x \neq 0,$$

$$\frac{H(t, \dot{x}(t), x(t))}{g(x(t))} = \frac{x^3(t)\dot{x}^2(t)}{x(t)t^3(\dot{x}^2(t)+1)(x^2(t)+1)} \leq \frac{1}{t^3} = m(t) \quad \forall x \neq 0, \dot{x} \in \mathfrak{R}, t > 0,$$

$$\lim_{t \rightarrow \infty} \frac{1}{t} \int_{t_0}^t \int_{t_0}^s [q(u) - m(u)] du ds = \lim_{t \rightarrow \infty} \frac{1}{t} \int_{t_0}^t \int_{t_0}^s \left[u^2 - \frac{1}{u^3} \right] du ds$$

$$\begin{aligned} &= \lim_{t \rightarrow \infty} \frac{1}{t} \int_{t_0}^t \left[\frac{s^3}{3} - \frac{1}{2s^2} - \frac{t_0^3}{3} + \frac{1}{2t_0^2} \right] ds \\ &= \lim_{t \rightarrow \infty} \frac{1}{t} \left[\frac{s^4}{12} + \frac{1}{2s} - \frac{t_0^3 s}{3} + \frac{s}{2t_0^2} \right]_{t_0}^t = \infty, \end{aligned}$$

It follows that by Theorem 2.1, Equation (10) is oscillatory.

Theorem 2.2: Suppose that the condition O_4 holds and

$$O_6 : \int_{t_0}^{\infty} \frac{ds}{r(s)} = \infty,$$

$$O_7 : \int_{t_0}^{\infty} [q(s) - m(s)] ds = \infty,$$

Then every solution of Equation (1) is oscillatory.

Proof. Let $x(t)$ be a non-oscillatory solution of Equation

(1), say $x(t) > 0$ for $t \geq T_0 \geq t_0$. Define

$$\omega(t) = \frac{r(t)\psi(x(t))\dot{x}(t)}{g(x(t))},$$

Then, for $t \geq T_0$, we get

$$\dot{\omega}(t) \leq m(t) - q(t),$$

Then, for every $t \geq T_0$, we obtain

$$\omega(t) \leq \omega(T_0) - \int_{T_0}^t [q(s) - m(s)] ds.$$

(11)

Then, by definition of ω , we get

$$\frac{r(t)\psi(x(t))\dot{x}(t)}{g(x(t))} \leq \omega(T_0) - \int_{T_0}^t [q(s) - m(s)] ds.$$

By condition O_7 there exists $T_1 \geq T_0$, such that

$$\dot{x}(t) < 0, \quad \text{for } t \geq T_1,$$

Also, by condition O_7 implies that there exists $T_2 \geq T_1$, such that

$$\int_{T_1}^{T_2} [q(s) - m(s)] ds = 0 \quad \text{and} \quad \int_{T_2}^t [q(s) - m(s)] ds \geq 0, \quad t \geq T_2.$$

Integrating Equation (1), we get

$$r(t)\psi(x(t))\dot{x}(t) = r(T_2)\psi(x(T_2))\dot{x}(T_2) + \int_{T_2}^t H(s, \dot{x}(s), x(s)) ds - \int_{T_2}^t G(s, x(s)) ds.$$

Then, for $t \geq T_2$, we have

$$\begin{aligned} r(t)\psi(x(t))\dot{x}(t) &\leq r(T_2)\psi(x(T_2))\dot{x}(T_2) - \int_{T_2}^t g(x(s))[q(s) - m(s)] ds \\ &\leq r(T_2)\psi(x(T_2))\dot{x}(T_2) - g(x(t)) \int_{T_2}^t [q(s) - m(s)] ds \\ &\quad + \int_{T_2}^t \dot{x}(s) g'(x(s)) \int_{T_2}^s [q(u) - m(u)] du ds. \end{aligned}$$

Hence, for $t \geq T_2$, we obtain

$$r(t)\psi(x(t))\dot{x}(t) \leq r(T_2)\psi(x(T_2))\dot{x}(T_2).$$

Then, for every $t \geq T_2$, we have

$$r(t)\dot{x}(t) \leq r(T_2)\psi(x(T_2))\dot{x}(T_2)l_2, \quad t \geq T_2.$$

Thus,

$$\dot{x}(t) \leq l_2 r(T_2)\psi(x(T_2))\dot{x}(T_2) \frac{1}{r(t)}.$$

(12)

Integrating the inequality (12) from T_2 to t , we obtain

$$x(t) \leq x(T_2) + l_2 r(T_2)\psi(x(T_2))\dot{x}(T_2) \int_{T_2}^t \frac{ds}{r(s)} \rightarrow -\infty \quad \text{as } t \rightarrow \infty,$$

which is a contradiction to the fact that $x(t) > 0$; hence, the proof is complete.

Example 2.2: Consider the equation:

$$\left[\left(\frac{t^2}{t^2+1} \right) \left(\frac{2x^2(t)+3}{2x^2(t)+4} \right) \dot{x}(t) \right]' + x(t) \left(\frac{1}{t} + x^4(t) \right) = \frac{x^3(t)\dot{x}^4(t)(\sin t - 1)}{(1+x^2(t))(\dot{x}^4(t)+1)}, \quad t > 0. \quad (13)$$

We note that

$$0 < r(t) = \frac{t^2}{t^2+1} < 1 \quad \text{and} \quad \int_{t_0}^{\infty} \frac{ds}{r(s)} = \int_{t_0}^{\infty} \left(1 + \frac{1}{s^2} \right) ds = \infty, \quad t > 0,$$



$$\psi(x) = \frac{2x^2 + 3}{2x^2 + 4} > 0 \quad \text{and} \quad \frac{1}{\psi(x)} = \frac{2x^2 + 4}{2x^2 + 3} > 1 \quad \text{for all } x \in \mathfrak{R},$$

$$\frac{G(t, x(t))}{g(x(t))} = \frac{x(t) \left(\frac{1}{t} + x^4(t) \right)}{x(t)} = \left(\frac{1}{t} + x^4(t) \right) \geq \frac{1}{t} = q(t) \quad \text{for all } x \neq 0 \text{ and } t \in [t_0, \infty)$$

and $xg(x) = x^2 > 0$ and $g'(x) = 1 > 0$ for $x \neq 0$,

$$\frac{H(t, \dot{x}(t), x(t))}{g(x(t))} = \frac{x^3 \dot{x}^4 (\sin t - 1)}{(1 + x^2)(\dot{x}^4 + 1)} \leq (\sin t - 1) = m(t) \quad \forall x, \dot{x} \in \mathfrak{R}, \text{ and } t \in [t_0, \infty)$$

$$\int_{t_0}^{\infty} [q(s) - m(s)] ds = \int_{t_0}^{\infty} \left[\frac{1}{s} - \sin s + 1 \right] ds = \ln s + \cos s + s \Big|_{t_0}^{\infty} = \infty,$$

It follows that by Theorem 2.2, Equation (13) is oscillatory.

Remark 2.1: Theorem 2.1 and Theorem 2.2 extend the results of Graef et. al. [6] and Li and Cheng [7].

3. References.

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بعض الشروط الكافية الجديدة لتذبذب المعادلات التفاضلية العادية الغير خطية

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